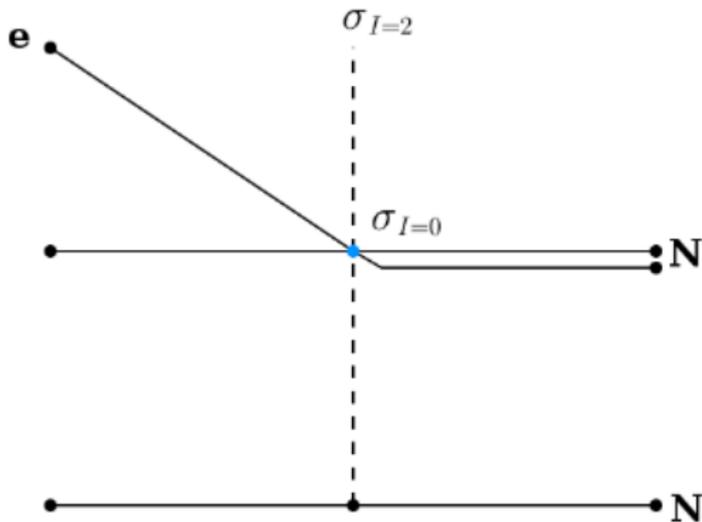


Working with theory about the Rossi effect



“Low Energy Nuclear Reactions”

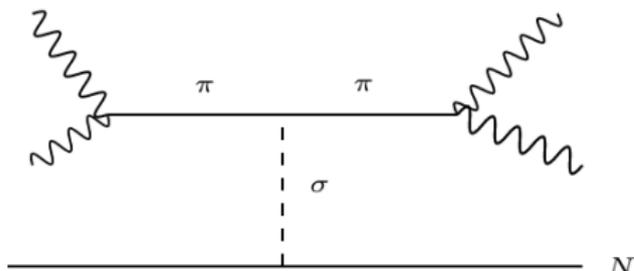
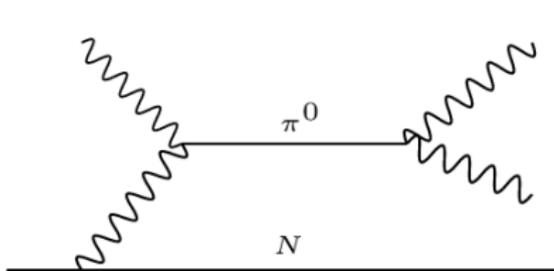
Problems to form a theory around:

- ▶ No strong radiation. If no detection of strong radiation is found together with evidence of isotopic shifts ie the strong force, then everything that creates strong radiation must be forbidden and a theory is formed around what is left.
- ▶ The limited range of the strong force.
- ▶ Nuclides don't come close enough at room temperature to affect each other with the strong force.

Solution presented here:

A new special potential of the strong force that is not found (common) in nature.

- ▶ Important feature is electron-nucleon interaction mediated by σ mesons.
- ▶ Releases energy continuously by slowly accelerate nuclides giving them kinetic energy.
- ▶ The special potential is a strong force potential triggered by electrons. Hence it does not require long range nucleon-nucleon interaction as a start point.



Outline

- ▶ Main theory in 3 steps
- ▶ Short on other theories
- ▶ Experiment
- ▶ Comparison theory to experiment
- ▶ Future

Outline

The main theory is developed in three steps:

1. Develop a strong force potential for nucleon-nucleon(N-N) interaction with the no γ radiation as a requirement.
 2. Enhance this potential by electron- σ meson interaction using isospin splitting of the σ meson in nucleons.
 3. Use nucleon polarizability theory to establish the electron nucleon interaction properties based on electromagnetic field component in the interaction.
- Note: Our paper¹ has theory presented in opposite order.

¹<https://arxiv.org/abs/1703.05249>

No γ problem for fusion

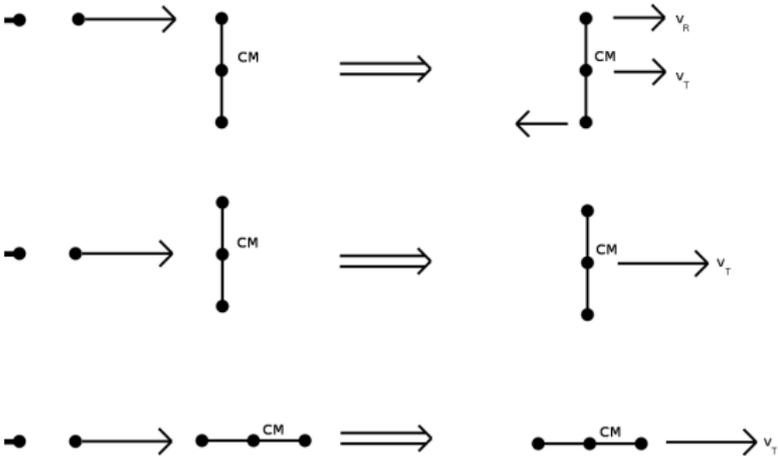
“cold fusion” is a bad word why? Because γ radiation needed for fusion reactions.

- ▶ Momentum and energy conservation can't be conserved in a pure $2 \rightarrow 1$ interaction (unless the sum of the initial momentum is 0).
- ▶ Simple problem in basic mechanics (non relativistic) momentum conservation proportional to v while energy conservation proportional to v^2 .
- ▶ Normal fusion has an extra photon to carry away the extra momentum.
- ▶ Nucleon transfer reaction is a $2 \rightarrow 2$ body reaction. Energy and momentum is allowed to be conserved without extra particles.

No excited state for nucleon transmission problem

Nucleon transmission reactions solves momentum conservation problem.

However just add a nucleon with a momentum transfer on a nuclide might create a oscillation motion unless the momentum transfer is applied on the center of mass.



No excited state for nucleon transmission problem

- ▶ Oscillation motion of nucleon inside nuclides=Excited states.
- ▶ Excited states in nuclides de-excite emitting strong radiation (most γ).
- ▶ Momentum transition to center of mass required for non excited state.

No excited state for nucleon transmission problem

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Potential from requirement:

- ▶ Nucleon transfer reaction needed.
- ▶ Momentum transfer must be applied on center of mass on both nuclides.
- ▶ Attraction to Mass=Scalar term needed in attractive potential.

Strong force

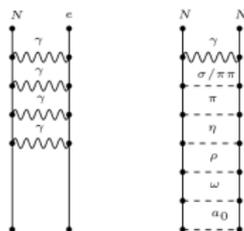
Energy range	fermion-boson (stable (\sim point-)particle -interaction particle)	Type	Main experimental need
$\sim 1\text{GeV} >$ keV-MeV	quark-gluon baryon(nucleon)-meson	Complete? Effective	decay of hadrons/mesons Nucleons in nuclides

- ▶ Mesons: 2 quark state separated by spin, charge, parity and quark generation(isospin: I, τ for generation 1).
- ▶ Baryons: 3 quark state Example:proton $p(I = 1/2, I_z = 1/2)$, neutron $n(I = 1/2, I_z = -1/2)$ and 4 Delta $\Delta(I = 3/2)$
- ▶ Effective theories uses LEC's(low energy constant) to describe phenomena in a lower energy range. Many theories exist depending on choice of approximative formula and problem.

Theoretical need for LENR:

- ▶ Nucleon-nucleon interaction.
- ▶ Nucleon electromagnetic interaction i.e. nucleon-electron/photon interaction.

Strong force



Nucleon nucleon(N-N) interaction:

- ▶ Using complete quark-gluon theory has problem with fermion doubling problem. Also quark theory can't explain the absent of electric dipole moment of the proton and neutron.
- ▶ nucleon nucleon interaction: Derived from meson exchange between nucleon. This is usually done as a trial function most succesfull theories fits constanst directly to r space operators.
- ▶ Both 2 and 3 nucleon interaction needed to explain all observations.

Strong force

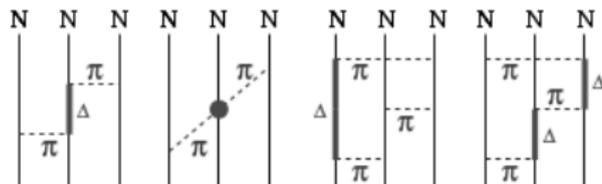
$$V_{L,S,I}(r) = \int \frac{d^3q}{(2\pi)^3} e^{iqr} \frac{g^2}{-q^2 - m^2} O(L, S, I) = -\frac{g^2}{4\pi} \frac{e^{-mr}}{r} O(L, S, I)$$

Meson	mass(MeV/c ²)	$I (J^P)$	role in N-N potential
π	138	$1 (0^-)$	Classic long range
σ	550	$0 (0^+)$	Binding
ω	782	$0 (1^-)$	repulsive
ρ	770	$1 (1^-)$	$L \cdot S$ interaction
a_0	980	$1 (0^+)$	short range
η	548	$0 (0^-)$	binding

Strong force

3 nucleon interaction:

- ▶ Internal structure change in the nucleon leads to different potential.
- ▶ Also the exchange particles interact in the space between two nucleons.



- ▶ Nucleon polarizability: Describes the internal changes of the nucleon when an external electromagnetic field is applied.

σ meson

- ▶ scalar meson, scalar term=center of mass i.e. property of nucleon potential derived from no γ requirement.
- ▶ In effective field theory: $\pi\pi$ s-wave resonance.
- ▶ Not natural long range. In one boson exchange potential(NN interaction usually change this term with one include a Δ):

$$V_{NN}^{(\sigma)}(r) = \int \frac{d^3q}{(2\pi)^3} e^{iqr} \frac{g_{\sigma NN}^2}{-q^2 - m_\sigma^2} = -\frac{g_{\sigma NN}^2}{4\pi} \frac{e^{-m_\sigma r}}{r}$$

with $m_\sigma \simeq 550$ MeV. Compare to EM potential from photon:

$$\frac{q}{4\pi} \frac{e^{-m_\gamma r}}{r}$$

with $m_\gamma = 0$.

σ meson isospin

σ interaction properties: The σ meson is a phase shift in $\pi\pi$ scattering separated by isospin states.

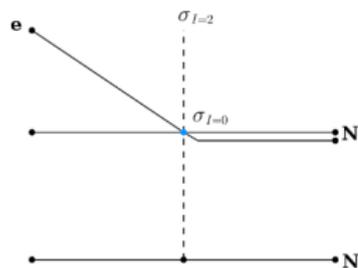
- ▶ From $\pi\pi$ scattering²: $m_{\sigma_{I=0}} = 36.77m_{\pi}^2$ and $m_{\sigma_{I=2}} = -21.62m_{\pi}^2$
- ▶ Hadron interaction properties: Proton and neutron isospin half state $\rightarrow m_{\sigma}$ in OBE potential is mixed:

$$m_{\sigma} = \sqrt{m_{\sigma_{I=0}}^2 - m_{\sigma_{I=2}}^2} \quad (550 \cong 543)$$

- ▶ Idea for step 2: electron-nucleon interaction enhance the range of the σ part of a N-N potential.

²G. Colangelo, J. Gasser, H. Leutwyler, $\pi\pi$ scattering, arxiv:hep-ph/0103088v1

σ meson isospin



- ▶ Why?: Electron is isospin 0 state interaction with $\sigma_{I=2}$ is suppressed compared to $\sigma_{I=0}$ which would increase the length of the σ part of the N-N potential.
- ▶ σ -electron interaction needs electron in nucleon since the interaction range of $\sigma_{I=0}$ is short.
- ▶ The σ -electron interaction doesn't equal a binding (by theory) therefore the electron must be kept in place by force.
- ▶ The electron must be able to spinflip to extract $\sigma_{I=0}$ energy out of the nucleon while stay in place.

Needed Electron nucleon interaction

- ▶ N-N interaction negligible on atomic level $\rightarrow e^-/\gamma - N$ interaction required.
- ▶ σ -electron interaction does not say anything about needed electron-nucleon interaction properties other than it requires electron at nucleon.
- ▶ If σ -electron interaction doesn't is attractive the distance between electron and nucleon must be in the 10^{-15} m range.
- ▶ electron nucleon interaction that would change the internal structure of the nucleon needed.

Nucleon polarizability

- ▶ What: Electromagnetic(EM) interaction of nucleon besides basic coulomb and magnet interaction i.e. internal structure changes.
- ▶ Compare to the 3 nucleon force where the internal structure changes affect the potential derived from the nucleon.
- ▶ Goal: Find EM interaction of e-N system that corresponds to binding condition of N-N force.
- ▶ Why? This would set the included nuclides in the new ground state binding condition after a nucleon transfer.
- ▶ Binding condition for particle systems are calculated by adding a negative term to the kinetic hamiltonian.
- ▶ Theoretical work have problems with choice of type of effective field approximation.

Nucleon polarizability

- ▶ Example coulomb interaction hamiltonian for atomic physics ($V < 0$ electromagnetic binding):

$$H\Psi = E\Psi \rightarrow (-k\nabla^2 + V)\Psi = E\Psi$$

- ▶ Polarizability calculated in perturbation theory by adding effective hamiltonians that are divided according to spacetime derivatives⁽ⁱ⁾ of the EM field

$$H = E_0 - \sum H_{\text{eff}}^{(i)}$$

- ▶ We look after conditions $H_{\text{eff}} > 0$.

Nucleon polarizability

Advance equations from perturbation theory³:

$$\begin{aligned} H_{\text{eff}}^{(2)} &= -\frac{1}{2}4\pi (\alpha_{E1}\bar{E}^2 + \beta_{M1}\bar{H}^2) \\ H_{\text{eff}}^{(3)} &= -\frac{1}{2}4\pi [\gamma_{E1E1}\bar{\sigma} \cdot (\bar{E} \times \dot{\bar{E}}) + \gamma_{M1M1}\bar{\sigma} \cdot (\bar{H} \times \dot{\bar{H}}) \\ &\quad - 2\gamma_{M1E2}E_{ij}\sigma_i H_j + 2\gamma_{E1M2}H_{ij}\sigma_i E_j] \\ H_{\text{eff}}^{(4)} &= -\frac{1}{2}4\pi (\alpha_{E1\nu}\dot{\bar{E}}^2 + \beta_{M1\nu}\dot{\bar{H}}^2) - \frac{1}{12}4\pi (\alpha_{E2}E_{ij}^2 + \beta_{M2}H_{ij}^2) \end{aligned} \quad (1)$$

$\alpha_x, \beta_x, \gamma_x$ = polarizability constants.

σ = Pauli spin matrices of the nucleon

E and H are components of the electromagnetic fields.

$$E_{ij} = \frac{1}{2}(\nabla_i E_j + \nabla_j E_i) \text{ (same for } H_{ij}\text{)}$$

Note: The third order perturbation is called spin polarizability and is not included in an classic static EM field.

³F.Hagelstein, R.Miskimen and V.Pascalutsa, "Nucleon Polarizabilities: from Compton Scattering to Hydrogen Atom," Prog. Part. Nucl. Phys. bf 88 (2016) 29 [arXiv:1512.03765 [nucl-th]].

Values

Theoretical and experimental values of the proton and neutron static dipole, quadrupole and dispersive polarizabilities. The units are $10^{-4} fm^3$ (dipole) and $10^{-4} fm^5$ quadrupole.

	α_{E1}	β_{M1}	α_{E2}	β_{M2}
Proton				
B χ PT Theory ⁴	11.2 ± 0.7	3.9 ± 0.7	17.3 ± 3.9	-15.5 ± 3.5
Experiment(PDG ⁵)	11.2 ± 0.4	2.5 ± 0.4		
Neutron				
B χ PT Theory	13.7 ± 3.1	4.6 ± 2.7	16.2 ± 3.7	-15.8 ± 3.6
Experiment(PDG)	11.8 ± 1.1	3.7 ± 1.2		

⁴V.Lensky and V.Pascalutsa, "Predictive powers of chiral perturbation theory in Compton scattering off protons," Eur. Phys. J. C 65 (2010) 195 [arXiv:0907.0451 [hep-ph]].

⁵C. Patrignani et al.(Particle Data Group), Chin. Phys. C, 40, 100001 (2016)

Values

Theoretical values of the proton and neutron static dispersive polarizabilities. The units are 10^{-4} fm^5 .

	$\alpha_{E1\nu}$	$\beta_{M1\nu}$
Proton		
B χ PT Theory	-1.3 ± 1.0	7.1 ± 2.5
Neutron		
B χ PT Theory	0.1 ± 1.0	7.2 ± 2.5

Values

Theoretical and experimental values of the proton and neutron static spin polarizabilities. The units are 10^{-4} fm^4 .

	γ_{E1E1}	γ_{M1M1}	γ_{E1M2}	γ_{M1E2}
Proton				
$B\chi$ PT Theory	-3.3 ± 0.8	2.9 ± 1.5	0.2 ± 0.2	1.1 ± 0.3
MAMI 2015 ⁶	-3.5 ± 1.2	3.16 ± 0.85	-0.7 ± 1.2	1.99 ± 0.29
Neutron				
$B\chi$ PT Theory	-4.7 ± 1.1	2.9 ± 1.5	0.2 ± 0.2	1.6 ± 0.4

► Sign of γ_{E1M2} visualize problem with different theories:

$O(p^4)_b$	$O(\epsilon^3)$	$O(p^4)_a$	K-Matrix	HDPV
0.7	1.0	0.2	-1.8	-0.02
DR	L_X	$HB\chi$ PT	$B\chi$ PT	MAMI 2015
-0.02	-0.7	-0.4 ± 0.4	-0.2 ± 0.2	-0.7 ± 1.2

⁶P.P.Martel et al. [A2 Collaboration], "Measurements of Double-Polarized Compton Scattering Asymmetries and Extraction of the Proton Spin Polarizabilities," Phys. Rev. Lett. 114 (2015) [arXiv:1408.1576 [nucl-ex]]

Polarizability binding conditions for electric field

- ▶ Define variable $x_{L,T} = \dot{\vec{E}}_{L,T} / \bar{E} \cdot \hat{\vec{E}} / \dot{\vec{E}}$ to get two differential equations:

$$\alpha_{E1} \pm \gamma_{E1E1} x_T + \alpha_{E1\nu} x_T^2 \quad (2)$$

The \pm sign is determined by the direction between the vectors $\bar{\sigma}$ and $(\bar{E} \times \dot{\vec{E}})$.

$$\alpha_{E1} + \alpha_{E1\nu} x_L^2 \quad (3)$$

Polarizability binding conditions for electric field

Calculations for theoretical values from $B\chi PT$ gives the $H_{eff} > 0$ ranges:

Nucleon	$\text{sgn} \bar{\sigma} \cdot (\bar{\mathbf{E}} \times \dot{\bar{\mathbf{E}}})$	$x = \dot{\bar{\mathbf{E}}}/\bar{\mathbf{E}}$ range (fm)
p	+	$x_T < -2$ $x_T > 4.5$
p	-	$x_T < -4.5$ $x_T > 2$
p	0	$x_L^2 > 0.11$
n	+	$3.1 < x_T < 44$
n	-	$-44 < x_T < -3.1$
n	0	-

Polarizability binding conditions for magnetic field

- ▶ $\beta_{M2} < 0$ gives $H_{eff} > 0$ at a center of magnetic quadrupole.

For combined electric and magnetic fields define $x = \sigma_i E_j / H_{ij}$ and E_j as $E \sin \theta$ (with θ the angle between dimension j and the plane defined by i and k). This gives the second order equation:

$$\beta_{M2}/6 + 2 \sin \theta \gamma_{E1M2} x + x^2 \alpha_1 \quad (4)$$

The condition for having $H_{eff} = 0$ values from the theoretical $B\chi PT$ values are fulfilled by:

$$\frac{6\gamma_{E1M2}^2}{\beta_{M2}} \sin \theta = \alpha_{E1} \quad (5)$$

$$H_{\text{eff}} > 0$$

The three conditions for $H_{\text{eff}} > 0$:

- ▶ A center of a magnetic quadrupole which also allows for a weak electric field.
- ▶ Two ranges from the parameter $\dot{\vec{E}}_{L,T}/\bar{E}$:
- ▶ The $x_L^2 > 0.11$ range has $\dot{\vec{E}}$ in the direction of \bar{E} .
- ▶ The x_T ranges has $\dot{\vec{E}}$ perpendicular to \bar{E} , this means a circular motion of the electron around the nucleon.

Atomic states

- ▶ Combine the short range need from the $\sigma_{I=0}$ mass with the spin polarizabilities yields special conditions on the nucleon electron relation.
- ▶ Due to the long range that the new potential is suppose to have, the electron has to have a stable position near the nuclide(in the fm range).
- ▶ Only atomic binding to have electron at nuclide is s-state atomic bindings however average distance is in the order of 10^{10} m.
- ▶ Solutions:
 1. Pressure would make the electron come closer to the nucleon.
 2. The nuclides creates a current that follows an electron current.

Atomic states

- ▶ The x_L solution can't be combined with the long time requirement since the interaction are a linear motion.
 - ▶ The magnetic quadrupole solution is compatible with both solution 1 and 2.
 - ▶ The x_T solution could be combined with solution 1 and 2 in special conditions:
1. Non s-state atomic binding would need extreme pressure to fulfill the range condition. For a s-state element the electron nuclide relation is approximately a rotation if the center of mass does not equal the center of charge. However the nucleon spin are in s-state perpendicular to the $(\vec{E} \times \dot{\vec{E}})$ vector not aligned. The solution is to have an electron that transform between a state with aligned nucleon spin and near nucleus condition. The right conditions are found for s- d_{z^2} overlaps.
 2. For nucleon current that follows electron current the nucleon spin must precise around the electron.

Other theories

New particle theories:

- ▶ Dark matter theories.
- ▶ Low energy virtual particles theory. I and Rossi are also making the hypothesis of the possibility that the temperatures of the plasma can reach the mass of , waves in fields that could annihilate without emitting high energy radiations because of the low energy.

Problem: Strong evidence of isotopic shifts requires link to the strong force at some point.

- ▶ Multi particle binding would explain the no γ condition with binding a lot of particles to the nucleon instead of one.

Problem: Electromagnet interaction strongly enhance one photon couplings.

Experiment

Observations⁷⁸⁹:

- ▶ Energy production without strong radiation.
- ▶ Isotopic shifts
- ▶ Positive ion current

⁷http://www.elforsk.se/Global/Omv%C3%A4rld_system/filer/LuganoReportSubmi

⁸K. A. Alabin, S. N. Andreev, A. G. Parkhomov. Results of Analyses of the Isotopic and Elemental Composition of Nickel- Hydrogen Fuel Reactors.

<https://drive.google.com/file/d/0B5Pc25a4cOM2cHBha0RLbUo5ZVU/view>

⁹<https://arxiv.org/abs/1703.05249>

No γ radiation

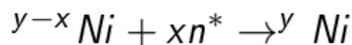
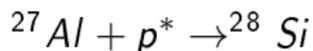
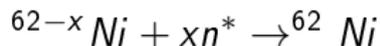
- ▶ Observed: Gamma radiation less than background level.
- ▶ For a detector ~ 0.5 m away this means $10^4 - 10^6$ γ /s for γ energies ~ 100 keV to 10 MeV.
- ▶ Observed: $10^{12} - 10^{15}$ transfer reactions/s. If each reaction creates ~ 1 MeV of energy.
- ▶ Creation of some radioactive nuclides that almost does not produce γ radiation still possible.
- ▶ Example:

${}^6\text{He}$: $\sim 10^{12}$ produced/s possible, above this secondary x-ray from β^- radiation should be detectable.

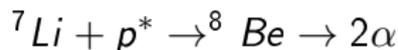
${}^{59}\text{Ni}$: $\sim 10^{20}$ produced/s possible. Above this rate 511 keV γ rays from positron annihilation should be at background level.

Isotopic shifts

Main detected isotopic shifts:



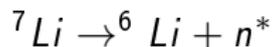
where p^* and n^* mean a bound nucleon. Also possible observed is:



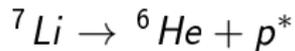
(If p is a free proton this would create measurable γ radiation above background level but not with bound)

Isotopic shifts

Neutron sources:



Proton sources:



Positive ion current

New experimental observation: Li/H ratio in plasma is related to output energy.

Output power is created when negative ions changes to positive ion kinetic energy.

Neutral plasma \rightarrow number and speed of positive and negative ions that enters the plasma are the same.

COP: Kinetic energy of positive ions/kinetic energy of negative ions.

Non relativistic kinetic energy:

$$\sum \frac{m_+ v_+^2}{2} / \sum \frac{m_- v_-^2}{2}$$

- ▶ Neutral plasma gives: $\sum v_+^2 = \sum v_-^2$
- ▶ COP is related to m_+/m_- i.e. in the range $m_{Li}/m_e = 14000$ to $m_H/m_e = 2000$.
- ▶ Measured COP in the doral test are in the range of thousands. Li/H ratio are reduced with the COP.

Important atomic states

Experimental observation of needed elements is in agreement with the theoretical requirement of atomic states.

- ▶ Free s-state electron elements needed to have spin flip electrons in nucleon.
- ▶ Free d_{z^2} electron elements needed to have nucleon spin perpendicular to electron.
- ▶ d_{z^2} -s overlap needed

d_{z^2} electron elements:

- ▶ Nickel group i.e. Nickel, Palladium, Platinum.

Free s-state electron elements:

- ▶ Hydrogen
- ▶ Alkalimetals: Lithium, Sodium, Potassium, Cesium.
- ▶ Some other metals: nickel, platinum, niobium, molybdenum, ruthenium, rhodium, and chromium.

Experiment-theory comparison summation

- ▶ No strong radiation \rightarrow continuous kinetic energy release of nuclides.
- ▶ (No γ) Momentum transfer is applied on center of mass \rightarrow potential is mediated by scalar meson (σ)
- ▶ Long range σ potential not natural. Created by special electron-nucleon interaction.
- ▶ One electron in s-state elements are needed: spin flip electron in the nucleon needed.
- ▶ One free electron in d_{z^2} shell elements needed: Tilt electron in right position compares to binding condition of polarizability.
- ▶ Plasma between Ni rods: Nickel creates $\sigma_{l=2}$ potential that drags protons and Lithium ions through air.

Future

Experimental to do:

- ▶ Important atomic states must be examined better. For example by doing isotopic shift measurement in slices.
- ▶ Find evidence for more possible isotopic shifts.
- ▶ Take α/β radiation spectrum to fit to theory.
- ▶ Exact numbers of output power compared to H/Li ratio in plasma.
- ▶ Hydrogen in plasma. Free hydrogen or proton in a long range nucleon transfer reaction?(would be possible to measure by the mass of the proton)

Not LENR experiment:

- ▶ Better polarizability measurement to confirm the theoretical values.

Theoretical to do:

- ▶ More exact theory for electron- σ interaction needed.
- ▶ Theory for multinucleus transfer reactions. Especially α clusters.(Deuteron-Palladium systems)
- ▶ Detailed study of available atomic states.